

# Cavity Ring-down Absorption Spectroscopy: Optical Characterization of ICI Product in Photodissociation of CH<sub>2</sub>ICI at 248 nm

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## Abstract

Iodine monochloride (ICl) elimination from one-photon dissociation of  $\text{CH}_2\text{ICl}$  at 248 nm is monitored by cavity ring-down absorption spectroscopy (CRDS). The spectrum of ICl is acquired in the transition of  $B^3\Pi_0 \leftarrow X^1\Sigma^+$  and is confirmed to result from a primary photodissociation; that is,  $\text{CH}_2\text{ICl} + h\nu \rightarrow \text{CH}_2 + \text{ICl}$ . The vibrational population ratio is determined with the aid of spectral simulation to be  $1:(0.36 \pm 0.10):(0.11 \pm 0.05)$  for the vibrational levels  $v = 0, 1$  and  $2$  in the ground electronic state, corresponding to a Boltzmann-like vibrational temperature of  $535 \pm 69$  K. The quantum yield of ICl molecular channel for the reaction is obtained to be  $0.052 \pm 0.026$  using a relative method in which the scheme  $\text{CH}_2\text{Br}_2 \rightarrow \text{CH}_2 + \text{Br}_2$  is adopted as the reference reaction. The ICl product contributed by the secondary collisions is minimized such that its quantum yield obtained is not overestimated. With the aid of CCSD(T)//B3LYP/MIDI! level of theory, the ICl elimination from  $\text{CH}_2\text{ICl}$  is evaluated to follow three pathways via either (1) a three-center transition state or (2) two isomerization transition states. But the three-center concerted mechanism is verified to be unfavorable.

## Abbreviations

CRDS, PES, RRKM, DFT

# 1 Introduction

Polyhalomethanes like  $\text{CH}_2\text{I}_2$ ,  $\text{CH}_2\text{Br}_2$ ,  $\text{CHBr}_3$ ,  $\text{CH}_2\text{BrI}$ ,  $\text{CH}_2\text{ClI}$ , and  $\text{CHBr}_2\text{Cl}$  are probably the main sources of reactive halogens (X) in the atmosphere. The subsequent reactions of reactive halogens with  $\text{O}_3$  to form XO radical initiate various tropospheric ozone depleting cycles.<sup>1</sup> Apart from ozone depletion, especially iodine plays a significant role in the formation of marine aerosol and cloud condensation nuclei, thus affecting the climate.<sup>2</sup> The main sources of polyhalomethanes are algae and phytoplankton present in ocean. Though these compounds dissolve in the oceanic water and marine boundary layer, measurable quantities of gases escape into the atmosphere.<sup>3</sup> Among the polyhalomethanes,  $\text{CH}_2\text{ICl}$  is one of the important and interesting compounds which contains both iodine and chlorine. For instance, various measurements<sup>4-6</sup> have shown seawater concentrations of  $\text{CH}_2\text{ICl}$  in the northwestern Atlantic to be comparable to or greater than that of  $\text{CH}_3\text{I}$ . Also it is reported that  $\text{CH}_2\text{ICl}$  is the potential source for atmospheric iodine compared to  $\text{CH}_3\text{I}$ . By solar irradiation, the  $\text{CH}_2\text{ICl}$  can be photo-dissociated to various products ( $\text{CH}_2$ ,  $\text{CH}_2\text{I}$ ,  $\text{CH}_2\text{Cl}$ ,  $\text{CHCl}$ ,  $\text{CHI}$ ,  $\text{I}$ ,  $\text{I}^*$ ,  $\text{Cl}$ ,  $\text{Cl}^*$ ,  $\text{HI}$ ,  $\text{HCl}$  and  $\text{ICl}$ ); such photo-dissociation may proceed both in solution and in the atmosphere.<sup>2</sup> In this study, the photodissociation of  $\text{CH}_2\text{ICl}$  was focused on the production of  $\text{ICl}$ ; this channel is optically characterized for the first time.

Various previous works have been dedicated to the studies of the photo-dissociation of  $\text{CH}_2\text{ICl}$  in gas phase. For example, Senapati et al.<sup>7,8</sup> used resonance-enhanced multiphoton ionization (REMPI) technique to study photo dissociation of  $\text{CH}_2\text{ICl}$  at different excitation wavelengths (222, 266, 280, and  $\sim 304$  nm). They measured the quantum yield for the production of  $\text{I}^*(^2\text{P}_{1/2})$  and  $\text{Cl}^*(^2\text{P}_{1/2})$  corresponding to the transitions from  $^3\text{Q}_0(\text{C-I})$  and  $^1\text{Q}_1(\text{C-I})$  states<sup>7</sup> and from the  $^3\text{Q}_0(\text{C-Cl})$  state,<sup>8</sup> respectively. The quantum yields of  $\text{I}^*$  and  $\text{Cl}^*$  obtained lie in between 0.22 and 0.55 at various photolysis wavelengths; among them,  $\text{I}^*$  yields a larger value for each wavelength.<sup>7,8</sup>

Tu et al.<sup>9</sup> used laser-induced fluorescence (LIF) technique to study the photolysis of different polyhalomethanes such as  $\text{CHI}_3$ ,  $\text{CH}_2\text{I}_2$ ,  $\text{CH}_3\text{I}$ , and  $\text{CH}_2\text{ICl}$  at 266 nm in a slow flow

cell at ambient temperature. Among the emission spectra of CH, C<sub>2</sub> and I<sub>2</sub> observed, they reported that most of the observed I<sub>2</sub> molecules result from the collisional product between the atomic iodine from photodissociation, instead of any molecular elimination channels.

Schmitt et al.<sup>10</sup> studied the photo decomposition of CH<sub>2</sub>ICl at 266 nm using gas chromatographic technique to analyze carbon containing products. They have identified eight products (C<sub>2</sub>H<sub>3</sub>Cl, C<sub>2</sub>H<sub>4</sub>, C<sub>2</sub>H<sub>2</sub>, C<sub>2</sub>H<sub>4</sub>Cl<sub>2</sub>, trans-C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub>, CH<sub>2</sub>Cl<sub>2</sub>, cis-C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub> and I<sub>2</sub>) which are formed mostly from secondary and recombination reactions. But they could not detect the primary photodissociation products such as HCl, HI and ICl. Other pathways were explored by Zhang et al.<sup>11</sup> using laser photo-fragmentation time-of-flight mass spectrometric method at 193 nm. They obtained the kinetic energy distributions for the predominant products such as I(<sup>2</sup>P<sub>1/2</sub>), I(<sup>2</sup>P<sub>3/2</sub>), Cl(<sup>2</sup>P<sub>1/2,3/2</sub>), singlet HCl and singlet ICl, and evaluated the branching ratio for the formation of I(<sup>2</sup>P<sub>1/2</sub>) and I(<sup>2</sup>P<sub>3/2</sub>). In addition, they theoretically predicted the production mechanism to CH<sub>2</sub> + ICl through an isomerization transition state. With an alternative technique of time-resolved electronic absorption spectroscopy, Preston et al.<sup>12</sup> probed the isomerization of CH<sub>2</sub>ICl into iso-CH<sub>2</sub>-Cl-I, which turned out to proceed in a high excited level with a brief lifetime prior to photo-dissociation.

Šulkova et al.<sup>13</sup> studied theoretically the reaction of OH radicals with CH<sub>2</sub>ICl, showing the atmospheric lifetime for the CH<sub>2</sub>ICl removal is about 1 year. They thus concluded that the atmospheric loss of CH<sub>2</sub>ICl due to the oxidation by OH radicals is not as rapid as with the oxidation by Cl atoms and photodissociation of CH<sub>2</sub>ICl by UV light.

Apart from the gas-phase reaction, a few studies were reported in photoexcitation of several polyhalomethanes including CH<sub>2</sub>ICl in dilute solutions using femtosecond transient absorption and time-resolved resonance Raman spectroscopy.<sup>14-18</sup> Åkesson and co-workers<sup>14</sup> reported the femtosecond transient absorption on the photolysis of CH<sub>2</sub>ICl in acetonitrile solution at room temperature. They observed iso-CH<sub>2</sub>Cl-I as first photoproduct with life time 100 ps and then ICl<sup>-</sup> as the second photoproduct at longer times. Philips group<sup>15</sup> investigated the initial formation and subsequent decay of the photoproduct produced following 267

nm excitation of CH<sub>2</sub>ICl in acetonitrile solution. They suggested that the iso-CH<sub>2</sub>Cl-I photoproduct is first produced from initially formed CH<sub>2</sub>ICl and I fragments via recombination within the solvent cage, followed by isomerization to give either the more stable iso-CH<sub>2</sub>Cl-I or parent CH<sub>2</sub>ICl molecules.

This work aims to study the ICl production channel using cavity ring-down absorption spectroscopy (CRDS) at 248 nm. A detailed characterization of this molecular pathway will be reported, such as (i) acquisition of ro-vibrational spectra of the ICl fragment in the  $B^3\Pi_0 \leftarrow X^1\Sigma^+$  transition, accompanied by the corresponding spectral simulation, (ii) determination of vibrational branching ratios of ICl ( $v''=0, 1, \text{ and } 2$ ) and the quantum yield for the ICl molecular channel, (iii) verification of ICl formation resulting from primary one-photon dissociation of CH<sub>2</sub>ICl (i.e.,  $\text{CH}_2\text{ICl} + h\nu \rightarrow \text{CH}_2 + \text{ICl}$ ), instead of a secondary reaction (e.g.,  $\text{I} + \text{CH}_2\text{ICl} \rightarrow \text{CH}_2\text{I} + \text{ICl}$ ), and (iv) ab initio calculations of dissociation pathways for the ICl channel, followed by evaluation of the Rice-Ramsperger-Kassel-Marcus<sup>19,20</sup> (RRKM) rate constants. Finally, the ICl production process is compared with Br<sub>2</sub>, and I<sub>2</sub> elimination from CH<sub>2</sub>Br<sub>2</sub> and CH<sub>2</sub>I<sub>2</sub>, respectively.

## 2 Experimental setup

The experimental setup of the CRDS apparatus has been described elsewhere<sup>21-27</sup> In brief, the CRDS is composed of three units, a ring-down flow cell, the photolysis-and-probe laser system and the ring-down time detection and signal analysis. A KrF excimer laser, emitting at 248 nm with the energy 10-30 mJ/pulse, was employed for photolysis. An Nd:YAG-pumped dye laser, operating with the dyes of sulforhodamine B (588-614 nm) or coumarin 503 (485-546 nm), was used as the ICl probe beam with the energy regulated about 2-4 mJ. The beam was guided through a spatial filter to select the TEM<sub>00</sub> mode.

Two sets of highly reflective mirrors were mounted to the long-arms of a cross-shaped stainless steel cavity, one for the 485-546 nm range and the other for the 588-614 nm range.

Both sets of mirrors have a reflectivity of 99.95% with 1 m radius of curvature and a diameter of 25.4 mm. The photolysis beam, focused via a 25-cm focal-length cylindrical lens, propagated through the short-arms to overlap with the CRDS probe beam perpendicularly. The ring-down time was monitored by a photo-multiplier tube sitting behind the rear mirror. To avoid pollution by chemical contact with the mirrors, an argon gas was kept to flow constantly across the cavity. Prior to introduction of the reagents, CH<sub>2</sub>ICl (Acros) and the ICl (Acros), a freeze-pump-thaw system at 77K was used up to 3-4 times. The absorption coefficient  $\alpha$  and the subsequent concentration of ICl may be determined by the fit of the acquired ring-down time ( $\tau$ ), according to the equation (1)

$$\alpha(\lambda) = \frac{d}{cl} \left[ \frac{1}{\tau} - \frac{1}{\tau_0} \right] \quad (1)$$

where  $d$  is the length of cavity,  $l$  is the length of light interacting with samples,  $c$  is the velocity of light, and  $\tau$  and  $\tau_0$  denote the ring-down time with and without the sample present in the ring-down cell, respectively. The ICl absorption spectra was recorded with a spectral resolution of 0.1 cm<sup>-1</sup>. The signal was acquired and analyzed with the aid of a MATLAB laboratory-developed software.

### 3 Theoretical methods

We characterized the adiabatic singlet ground-state potential energy surface (PES) of the various dissociation channels of CH<sub>2</sub>ICl dissociation. The geometries of the species involved (reactant, transition states and products) in the CH<sub>2</sub>ICl dissociation were optimized using hybrid density B3LYP<sup>28,29</sup>/MIDI!<sup>30,31</sup> functional theory. Harmonic frequencies and the zero point energies of all the species were obtained using the same level of theory. Both reactant and products were obtained with all positive frequencies and the transition state with one imaginary frequency. We further performed the intrinsic reaction coordinate (IRC) using the Gonzalez and Schlegel<sup>32,33</sup> second-order algorithm to verify the connections between the

transition state and the reactant and products. Finally, we employed the CCSD(T)/MIDI! method on optimized geometries to get a more accurate energy. The basis set MIDI! was utilized for its moderate size for the heavy atom I, and succeeded in predicting a reasonable ICl bond length. These calculations were carried out using the Gaussian09 suite of program.<sup>34</sup> The RRKM rate constant was calculated for the formation of ICl from photo-dissociation of ground state of CH<sub>2</sub>ICl at 248nm.

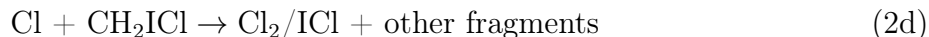
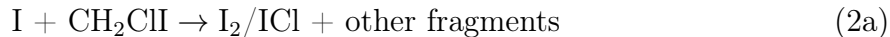
## 4 Results and Discussion

### 4.1 Rovibrational spectra of ICl fragment

Using CRDS, we acquired the ro-vibrational spectra of ICl ( $v'' = 0, 1$  and  $2$ ) from 591 to 606 nm in the  $B^3\Pi_0 \leftarrow X^1\Sigma^+$  transition produced by the photo-dissociation of the CH<sub>2</sub>ICl at 248 nm. Fig. 1 shows a portion of the CRDS spectrum of the ICl fragment, accompanied by the pure ICl spectrum in the same region. The agreement between the spectra is reasonably good and leads to the conclusion that ICl is produced from the CH<sub>2</sub>ICl photo-dissociation and the spectral region is not interfered by other fragmented species.

To check if there are any secondary reactions which contributed to ICl, we compare with the pure I<sub>2</sub> spectrum recorded under otherwise the same conditions. As shown in Fig. 2, no significant trace of I<sub>2</sub> is found in the ICl spectrum. As a result, one can conclude that the ICl fragment is eliminated by the primary photo-dissociation of CH<sub>2</sub>ICl and its contribution by the secondary reaction should be negligible. If the secondary reaction does happen, one might observe formation of either product of ICl, Cl<sub>2</sub> and I<sub>2</sub> according to the following

secondary reactions (2a-2f).



Among the above products, a larger amount of  $\text{I}_2$  is expected in the secondary reactions, and a larger quantum yield of  $\text{I}^*$  with respect to  $\text{Cl}^*$  was also detected in the photolysis at varied wavelengths,<sup>7,8</sup> because the C-I bond energy is smaller than the C-Cl bond energy. In addition, the experimental condition of  $\text{CH}_2\text{ICl}$  at 34 mTorr (Fig. 5) corresponds to a concentration of  $1.2 \times 10^{15}$  molecules/ $\text{cm}^3$ . The I atomic elimination should be the major channel, which is assumed to take a yield of 50 %, corresponding to  $6 \times 10^{14}$  molecules/ $\text{cm}^3$ . Given a relative speed  $3 \times 10^4$  cm/s at 300 K between  $\text{CH}_2\text{ICl}$  and I and the collision cross section of  $35 \text{ \AA}^2$ , each  $\text{CH}_2\text{ICl}$  molecule may encounter  $\sim 0.3$  collisions by the released I during the ring-down time of 5  $\mu\text{s}$ . Thus, the ICl product resulting from the secondary collision (i.e.,  $\text{I} + \text{CH}_2\text{ICl} \rightarrow \text{ICl} + \text{CH}_2\text{I}$ ) is negligible. The calculated result is consistent with the experimental finding. In addition, the root-mean-squared distance that the ICl molecules diffuse within a detection period of 4-5  $\mu\text{s}$  ring-down time is about 1 mm, which is confined within the cross section (diameter = 3 mm) along the detection axis and within the photolysis-probe delay time of 20 ns.

To find out how many photon number are involved in the photolysis, rotational line intensity of ICl spectrum at 519.685 nm was measured as a function of the photolysis laser energy. As shown in Fig. 3b, the resulting plot yields a straight line which supports the



single-photon involvement in the molecular elimination of ICl. In addition to the above measurements, the CH<sub>2</sub>ICl pressure dependence of the ICl product was conducted. Given a constant photolysis laser energy, the plot of the ICl concentration as a function of the CH<sub>2</sub>ICl pressure (0.1 -2 Torr) yields a straight line, as shown in Fig 4. Because the channel of halogen atomic elimination dominates over the dissociation pathways, the pseudo-first order reaction cannot be applied to the secondary reactions (2a-2f). Thus, the ICl product contributed from the reaction of a halogen atom with precursor requires two CH<sub>2</sub>ICl molecules involved. The result lends the support to the fact that the ICl contribution from the secondary reactions should be minimal.

The CRDS spectral regions were carefully selected to contain the  $v'' = 0, 1$  and  $2$  bands. For instance, the ICl spectrum in the 595.80-596.60 nm region contains the vibrational bands  $(v', v'') = (14,0)$  and  $(13,1)$  of the  $B^3\Pi_0 \leftarrow X^1\Sigma^+$  transition. Similarly, the spectrum in the region of 598.0-598.4 nm corresponds to the bands of  $(v', v'') = (14,0), (16,1)$  and  $(1,2)$ . Fig.5 displays the ICl fragment in the 598.0-598.4 nm region, in which we also include the pure ICl spectrum and the simulated counterpart for individual vibrational band with appropriate adjustment of the population ratio of  $\text{ICl}(v''=1)/\text{ICl}(v''=0)$  and  $\text{ICl}(v''=2)/\text{ICl}(v''=1)$ . When  $\text{ICl}(v''=1)/\text{ICl}(v''=0)$  is fixed at 0.36 and  $\text{ICl}(v''=2)/\text{ICl}(v''=1)$  is adjusted to 0.11, the spectral simulation becomes appreciably consistent with the experimental finding.

Given the same population ratio for ICl ( $v''=0,1$ , and  $2$ ), Fig. 6 shows comparison of ICl fragment with the simulated counterpart in the region of 605-605.8 nm which contains the  $(v', v'')$  bands of  $(14,0), (13,1), (14,1), (0,2)$  and  $(2,3)$ . The population of  $v''=3$  is buried seriously in those populations of  $v''=0, 1$ , and  $2$  and thus can not be evaluated accurately.

## 4.2 Simulation and rotational populations

Despite a long ring-down time averaged to 5  $\mu\text{s}$ , the vibrational population may be preserved approximately in the nascent state for the vibrational energy transfer is slow, but the rotational population is rapidly thermally equilibrated. The reported Franck-Condon factors

(FCFs) are available in the literature only for a few transitions.<sup>35,36</sup> Thus, it is required to calculate the potential energy curves involved in the transition and subsequently evaluate the FCFs for simulation of the ICl spectra. To construct these potential energy curves, the spectroscopic parameters of  $B^3\Pi_0$  and  $X^1\Sigma^+$  electronic states are obtained from the previous work.<sup>37</sup> Then, the potential energy curves for the  $B^3\Pi_0$  and  $X^1\Sigma^+$  electronic states of ICl can be calculated using Rydberg<sup>38,39</sup>-Klein<sup>40</sup>-Rees<sup>41</sup> (RKR) method.<sup>42</sup> The FCFs was evaluated as the extent of overlap integral between the vibrational wavefunctions of the  $B^3\Pi_0$  and  $X^1\Sigma^+$  electronic states with the aid of LEVEL program.<sup>43</sup> Given FCFs obtained (see Supporting Information, SI), then the spectral intensities of ICl spectrum can be calculated using the equation (3)<sup>44</sup>

$$I = k \frac{(FCF)(HLF)}{2J'' + 1} N_{J''} r_{iso} f \quad (3)$$

where  $k$  is a factor associated with instrument and experimental conditions,  $FCF$  the Franck-Condon factors,  $HLF$  the Hönl-London-London factor,  $J''$  the rotational quantum number of the ground state ICl,  $N_{J''}$  the Boltzmann distribution of the rotational population for which 300 K is assumed,  $r_{iso}$  the ratio of isotopic variants,  $I^{35}\text{Cl}:I^{37}\text{Cl}$  equal to 0.7576 : 0.2424, and  $f$  the intensity ratio of strong and weak rotational lines.

The molecular constants of ground ( $X^1\Sigma^+$ ) and excited ( $B^3\Pi_0$ ) states of  $I^{35}\text{Cl}$  used in the present study are given in Table 1. As for  $I^{37}\text{Cl}$ , its molecular constants were derived by applying isotopic effect.

Table 1: molecular constants ( $\text{cm}^{-1}$ ) of ground ( $X^1\Sigma^+$ ) and excited ( $B^3\Pi_0$ ) states of  $I^{35}\text{Cl}$ .

Electronic state	$w_e(\text{cm}^{-1})$	$w_e x_e(\text{cm}^{-1})$	$w_e y_e(\text{cm}^{-1})$	$D_e(\text{cm}^{-1})$	$B_e(\text{cm}^{-1})$
$X^1\Sigma^+$	384.293	1.501		$4.03 \times 10^{-8}$	0.0872
$B^3\Pi_0$	204.5	2.595	-0.937	$1.00 \times 10^{-7}$	0.0872

To estimate the uncertainty for the simulation, a comparison between the experimental spectra and the simulation evaluated for different population ratios in the regions of 596.50-596.55 nm and 598.09-598.12 nm is shown in the Fig. 7. A population ratio  $\text{ICl}(v'' =$

0)/ICl( $v'' = 1$ ) is then optimized to be  $0.36 \pm 0.10$  in the range of 596.50-596.55 nm. Given ICl( $v'' = 0$ )/ICl( $v'' = 1$ ) fixed at 0.36, the ICl( $v'' = 0$ )/ICl( $v'' = 2$ ) ratio in the 598.09-598.12 nm range can be optimally adjusted to  $0.11 \pm 0.05$ . Accordingly, the population ratio of  $v'' = 0:1:2$  is equal to  $1:(0.36 \pm 0.10):(0.11 \pm 0.05)$  corresponding to the Boltzmann vibrational temperature  $535 \pm 69$  K .

### 4.3 Quantum yield for ICl channel

The quantum yield for the formation of ICl from the CH<sub>2</sub>ICl photo-dissociation is evaluated by  $[ICl]/N_p$ , in which  $[ICl]$  denotes the ICl concentration produced in the laser beam-crossed region and  $N_p$  the number density of photon absorbed in the same region. Alternatively, the quantum yield for the ICl channel can be determined by the indirect method, when the quantum yields of similar reactions are well established. The Br<sub>2</sub> formation from CH<sub>2</sub>Br<sub>2</sub> photo-dissociation studied in a previous work was chosen as reference reaction with the quantum yield of  $0.21 \pm 0.06$ <sup>21,27</sup>. The following equation is used to calculate the quantum yield of ICl production

$$\frac{\Phi_{ICl}}{\Phi_{Br_2}} = \frac{[ICl]n_{CH_2Br_2}\sigma_{CH_2Br_2}}{[Br_2]n_{CH_2ICl}\sigma_{CH_2ICl}} \quad (4)$$

where  $[ICl]$  and  $[Br_2]$  are concentrations obtained from the experiment,  $n_{CH_2ICl}$  and  $n_{CH_2Br_2}$  the number density of CH<sub>2</sub>ICl and CH<sub>2</sub>Br<sub>2</sub>,  $\sigma_{CH_2ICl}$  and  $\sigma_{CH_2Br_2}$  are the absorption cross sections of CH<sub>2</sub>ICl and CH<sub>2</sub>Br<sub>2</sub> at 248 nm. We evaluated Br<sub>2</sub> concentration to be  $8.1 \times 10^{15}$  molecules/cm<sup>3</sup> at the rotational line of 519.68 nm by the ring-down time measurements, and  $n_{CH_2Br_2}$  was experimentally controlled to 596 mTorr. By analogy, the ICl concentration was evaluated to be  $4.7 \times 10^{14}$  molecules/cm<sup>3</sup> at the almost identical rotational line, and  $n_{CH_2ICl}$  was controlled at 300 mTorr. Given  $\sigma_{CH_2Br_2}$  evaluated to be  $3.7 \times 10^{-19}$  cm<sup>2</sup>,  $\sigma_{CH_2ICl}$  equal to  $1.52 \times 10^{-19}$  cm<sup>2</sup> and the Br<sub>2</sub> quantum yield of  $0.2 \pm 0.06$ , substituting the above values into the equation yields a quantum yield of  $0.052 \pm 0.026$  for ICl elimination

from  $\text{CH}_2\text{ICl}$  photolysis. The uncertainty in the quantum yield value was estimated using error propagation formula as given in equation 5.

$$\Delta\phi_{\text{ICl}} = \phi_{\text{ICl}} \sqrt{\left(\frac{\Delta[\text{ICl}]}{[\text{ICl}]}\right)^2 + \left(\frac{\Delta[\text{Br}_2]}{[\text{Br}_2]}\right)^2 + \left(\frac{\Delta\sigma_{\text{CH}_2\text{Br}_2}}{\sigma_{\text{CH}_2\text{Br}_2}}\right)^2 + \left(\frac{\Delta\sigma_{\text{CH}_2\text{ICl}}}{\sigma_{\text{CH}_2\text{ICl}}}\right)^2 + \left(\frac{\Delta n_{\text{CH}_2\text{Br}_2}}{n_{\text{CH}_2\text{Br}_2}}\right)^2 + \left(\frac{\Delta n_{\text{CH}_2\text{ICl}}}{n_{\text{CH}_2\text{ICl}}}\right)^2 + \left(\frac{\Delta\phi_{\text{Br}_2}}{\phi_{\text{Br}_2}}\right)^2}$$

(5)

In comparison with this result, the quantum yields of  $\text{Br}_2$  and  $\text{I}_2$  produced from photodissociation of  $\text{CH}_2\text{Br}_2$ ,<sup>23,27</sup> and  $\text{CH}_2\text{I}_2$ ,<sup>26</sup> at the same wavelength yielded the values of  $0.21 \pm 0.06$  and  $0.004 \pm 0.0025$  respectively. The quantum yield for the  $\text{ICl}$  channel from  $\text{CH}_2\text{ICl}$  is smaller by a factor of four than that of  $\text{Br}_2$  from  $\text{CH}_2\text{Br}_2$ , but ten times larger than  $\text{I}_2$  from  $\text{CH}_2\text{I}_2$ . The obtained order of quantum yields of  $\text{Br}_2$ ,  $\text{ICl}$  and  $\text{I}_2$  apparently does not follow the strength of the C-X bond rupture. It is impossible to interpret the trend without taking into account the branching ratios of X elimination in the excited and ground state of halo-carbons and the efficiency of internal conversion from the state  $B^3\Pi_0$  to  $X^1\Sigma^+$ .

#### 4.4 Photodissociation mechanism

The density functional theory was applied to study the photo dissociation mechanism of the  $\text{CH}_2\text{ICl}$  molecule. The analysis of the PES was carried out and the geometries for reactant ( $\text{CH}_2\text{ICl}$ ), the transition states (tsICl, ts $\text{CH}_2\text{-I-Cl}$ , ts $\text{CH}_2\text{-Cl-I}$ , tsCHI and ts CHCl) and products ( $\text{CH}_2$ , ICl,  $\text{CH}_2\text{-I-Cl}$ , HI, HCl, CHI and CHCl) were optimized by the aid of the B3LYP/MIDI! level of theory. The calculations lead to a result of no negative frequencies for the  $\text{CH}_2\text{ICl}$ ,  $\text{CH}_2$ , ICl,  $\text{CH}_2\text{-I-Cl}$ , HI, HCl, CHI and CHCl, but a negative frequency for all transition states. To confirm that the transition states are connected with the reactant and the products, IRC calculations were conducted on all the transition states for the different reactions paths.

The high level CCSD(T) calculations were further carried out on the B3LYP optimized ge-

ometry to obtain more accurate energy. The B3LYP/MIDI! zero-point and CCSD(T)/MIDI! energies of all geometries are tabulated in S.I. The optimized geometries for the three-center elimination path and the two possible isomerization paths are displayed in Fig. 8 with inter-nal coordinates. As shown in the figure the bond distances for the C-I, C-Cl in the reactant are up to 1.802 Å and 2.188 Å respectively. In contrast, the corresponding bond distances are elongated to 2.273 Å and 3.005 Å in the transition state tsCH<sub>2</sub>ICl. In the case of the ICl product, the equilibrium bond length of Cl-I is 2.336 Å while the corresponding distance in the transition state is about 2.430 Å. The bond angle of I-C-Cl is shortened to 68.5° in the tsCH<sub>2</sub>ICl from 114° in the reactant.

The ground state CH<sub>2</sub>ICl decomposes into CH<sub>2</sub> and ICl via tsCH<sub>2</sub>ICl with a barrier height of 437 kJ/mol. The reaction scheme is endothermic with the energy of 436 kJ/mol. The dissociation pathways for the CH<sub>2</sub>ICl to form ICl were evaluated by the theory of CCSD(T)//B3LYP/MIDI! level and shown in Fig. 9.

In addition to PES exploration, density and sum of states were calculated for reactant and transition states using Baeyer-swineheart algorithm implemented in multiwell program.<sup>45,46</sup> The unimolecular rate constants for the dissociation channel (CH<sub>2</sub>ICl → CH<sub>2</sub> + ICl) were then calculated using equation (6).

$$k(E) = \frac{1}{h} \frac{G(E)}{N(E - E_r)} \quad (6)$$

where  $h$  is the Planck’s constant,  $G(E)$  is the sum of states for the active degrees of freedom at the transition state and  $N(E - E_r)$  is the vibrational density of states for the active degrees of freedom in the reactant. The calculated sum of vibrational states for the different path’s transition states as a function of the excitation energy are shown in Fig. S1. Also the density of states for the reactant (CH<sub>2</sub>ICl) as a function of the excitation energy are shown in Fig. S2. Comparison of the rate constants between three-center elimination and the isomerization pathway is given in Table 2. Despite much larger rates to surpass

the isomerization transition states ( $k_2$  and  $k_3$  paths), the net rates are impeded by the endothermicity. Similar pathway with isomerization transition state is also found in the elimination of  $I_2$  from  $CH_2I_2$ .<sup>28</sup>

Zhang et al.<sup>11</sup> have calculated the dissociation pathway in which the  $ICl$ (singlet state) +  $CH_2$ (singlet state) were produced only via an isomerization transition state,  $Cl$  sitting in between  $C$  and  $I$  (the  $k_3$  path). In contrast, we find out three reaction mechanisms, one with three-center concerted elimination and the other two surpassing isomerization transition state ( $k_2$  and  $k_3$  paths in Fig. 9). Despite the photolysis energy higher than the barrier along the  $k_1$  path, the three-center concerted elimination should result in a triplet  $CH_2$  and triplet  $ICl$  for electron-spin conservation. The ground triplet state  $CH_2$  has been verified unfavorable in kinetic energy analysis of the  $CH_2 + ICl$  channel.<sup>11</sup> In addition, the  $ICl$  fragment detected with CRDS is confirmed to lie in the singlet ground state. Thus, it is difficult for the  $CH_2ICl$  molecule to follow the  $k_1$  path leading to the  $ICl$  fragment. Apart from the  $k_3$  path reported previously by Zhang et al., we find herein that the  $k_2$  path proceeding via an isomerization transition state with  $I$  sitting in between  $C$  and  $Cl$  is available to produce the  $CH_2 + ICl$  channel.

In the gas reaction,  $ICH_2Cl$  proceeds to  $CH_2 + ICl$  via an isomerization transition state of  $CH_2I-Cl$  or  $CH_2Cl-I$  with the barrier of 352 or 282 kJ/mol, respectively. In contrast, the isomerization step  $CH_2I-Cl$  with a much lower energy barrier of 167 kJ/mol was also found in the solution phase via a recombination between  $CHI$  and  $Cl$  by the solvent cage effect. Despite the same isomerization step, the reaction mechanism between gas and liquid phases turn out to be different.

Table 2: Rate constants comparison between different paths.

Reactions	Rate constants ( $s^{-1}$ ) at 248 nm
$CH_2ICl \rightarrow CH_2 + ICl$	$k_1 = 1.026 \times 10^8$
$CH_2ICl \rightarrow CH_2-I-Cl$	$k_2 = 1.451 \times 10^{10}$
$CH_2ICl \rightarrow CH_2-Cl-I$	$k_3 = 1.637 \times 10^{11}$

## 5 Conclusions

The ICl fragment elimination from the photolysis of  $\text{CH}_2\text{ICl}$  at 248nm was demonstrated by using the CRDS technique. The acquired ICl ( $v''=0, 1, \text{ and } 2$ ) spectra in the  $B^3\Pi_0 \leftarrow X^1\Sigma^+$  transition are simulated to obtain the vibrational population ratio and the corresponding Boltzmann vibrational temperature. The quantum yield of the ICl channel was estimated. Its value lies in between the  $\text{Br}_2$  and  $\text{I}_2$  elimination from  $\text{CH}_2\text{Br}_2$  and  $\text{CH}_2\text{I}_2$ , respectively. However, interpretation for the trend of the quantum yields cannot simply rely on the strength of the C-X bond rupture. We also verify that the ICl fragment should result from the reaction,  $\text{CH}_2\text{ICl} + h\nu \rightarrow \text{CH}_2 + \text{ICl}$ , involving one-photon mechanism, and the ICl contribution from secondary reactions is negligible. With the aid of theoretical calculations, three dissociation pathways are calculated. The ICl product can be formed via either a three-center concerted transition state or two isomerization transition states. While taking into account the electron-spin conservation, the latter path following isomerization mechanism is favored.

## 6 Supporting Information

The Franck-Condon factors correspond to the ground ( $X^1\Sigma^+$ ) and excited ( $B^3\Pi_0$ ) states, sum and density of states correspond to optimized reactant and transition states, Energies of various species involved in the photo-dissociation of  $\text{CH}_2\text{ICl}$  obtained at B3LYP and CCSD(T) level of theories.

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for High-performance Computer of Taiwan were utilized in the calculations.

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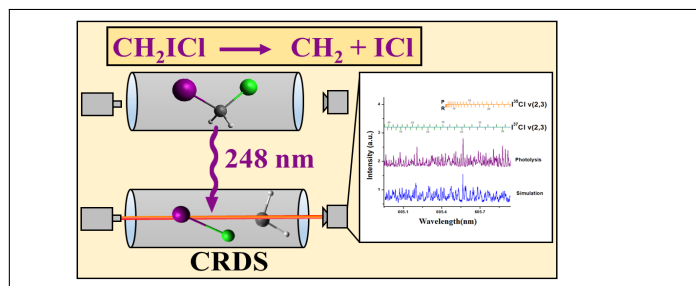
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# Graphical TOC Entry



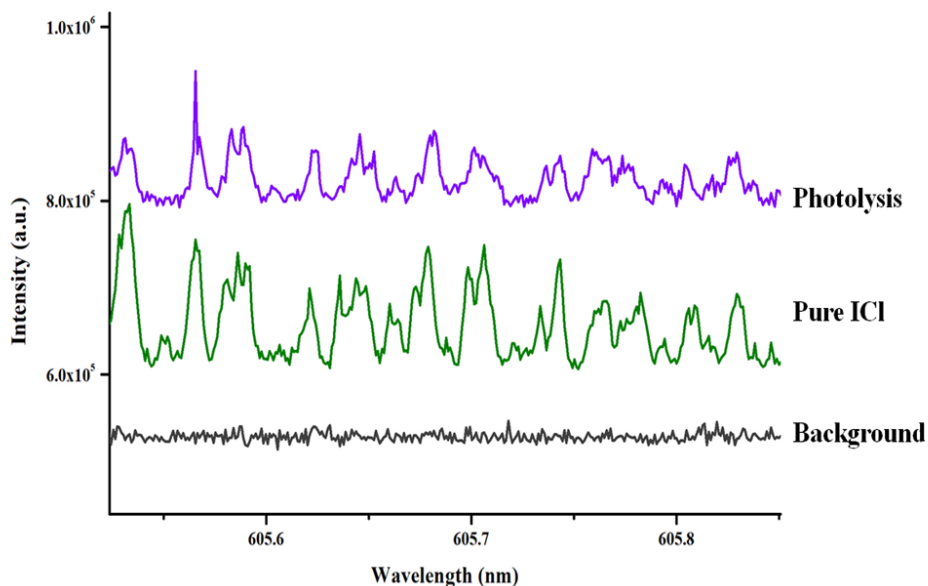


Figure 1: Comparison between a portion of ICl spectra acquired in the photolysis of  $\text{CH}_2\text{ICl}$  at 248 nm ) at 0.006 Torr (violet line and pure ICl molecule (green line) and the background spectrum) at 0.034 Torr (grey line). Both spectra are taking at room temperature.

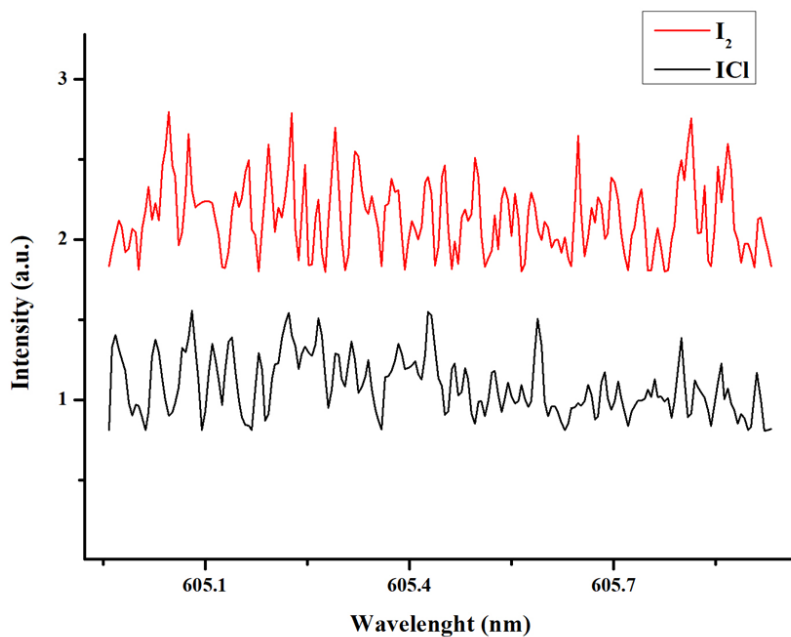


Figure 2: Comparison between a portion of pure ICl spectra (red line) and pure  $\text{I}_2$  molecule at room pressure (black line). Both spectra are taking at room temperature.

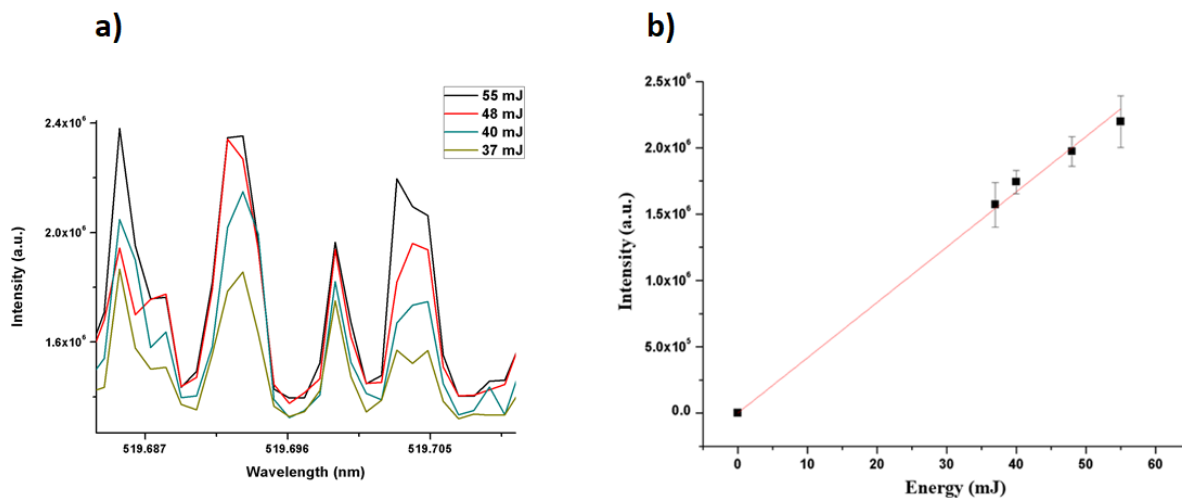


Figure 3: Intensity of ICl fragment as a function of photolysis laser energy varied from 37 to 55 mJ with a pressure of 0.021 Torr at room temperature. The rotational line at 519.685 nm is selected for this measurement.

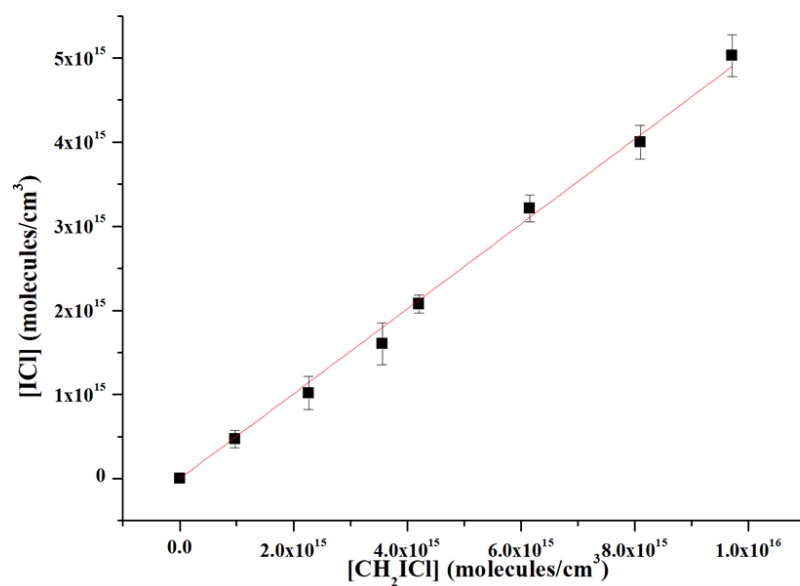


Figure 4: The amount of ICl fragment produced as a function of CH<sub>2</sub>ICl concentration varied from  $1 \times 10^{15}$  to  $8 \times 10^{15}$  molecules cm<sup>-3</sup> at room temperature at 248nm using a laser power of 39 mJ. The rotational line at 519.685 nm is selected for this measurement.



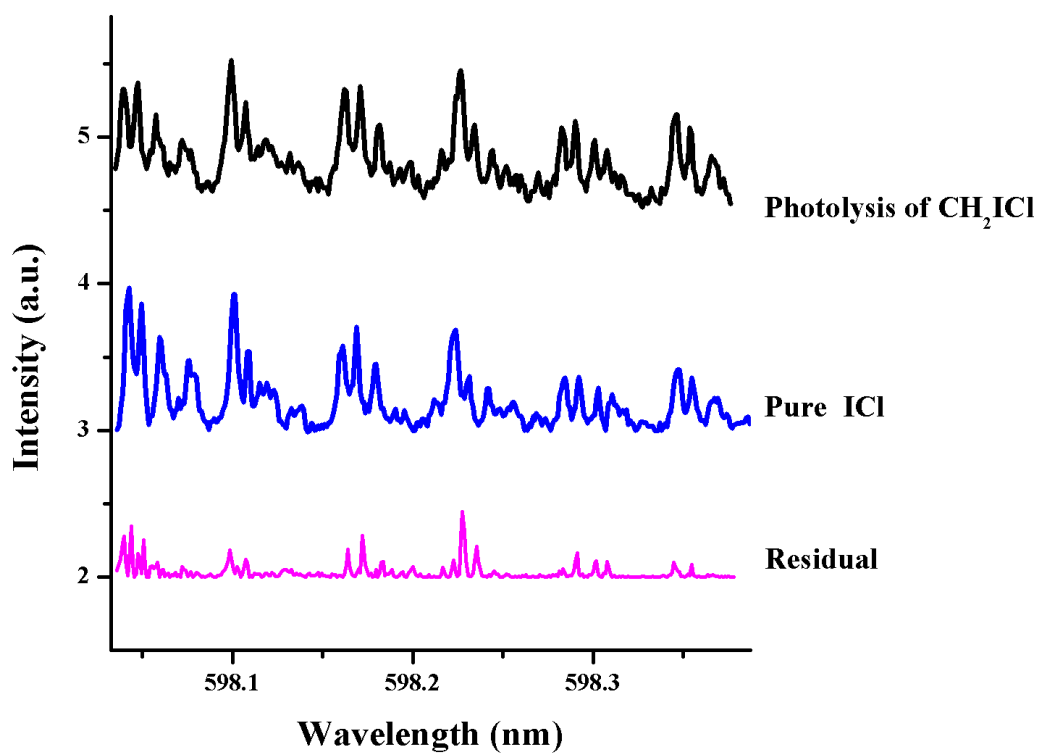


Figure 5: A portion of ICl spectra acquired in the photolysis of  $\text{CH}_2\text{ICl}$  at 248 nm. ICl spectrum from photolysis of  $\text{CH}_2\text{ICl}$  at 248 nm at 0.034 Torr (black line), pure ICl molecule at 0.086 Torr (blue line), and residual spectrum (pink line). Both spectra are taken at room temperature.

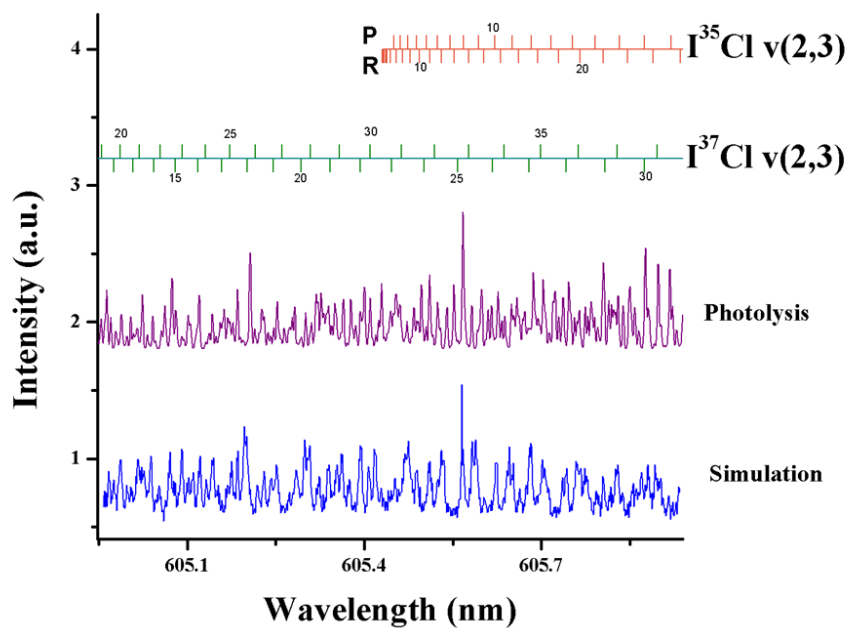


Figure 6: A portion of ICl spectra acquired in the photolysis of  $\text{CH}_2\text{ICl}$  at 248 nm. ICl spectrum from photolysis of  $\text{CH}_2\text{ICl}$  at 248 nm at 0.044 Torr and room temperature (violet line) and simulation of ICl molecule (blue line)

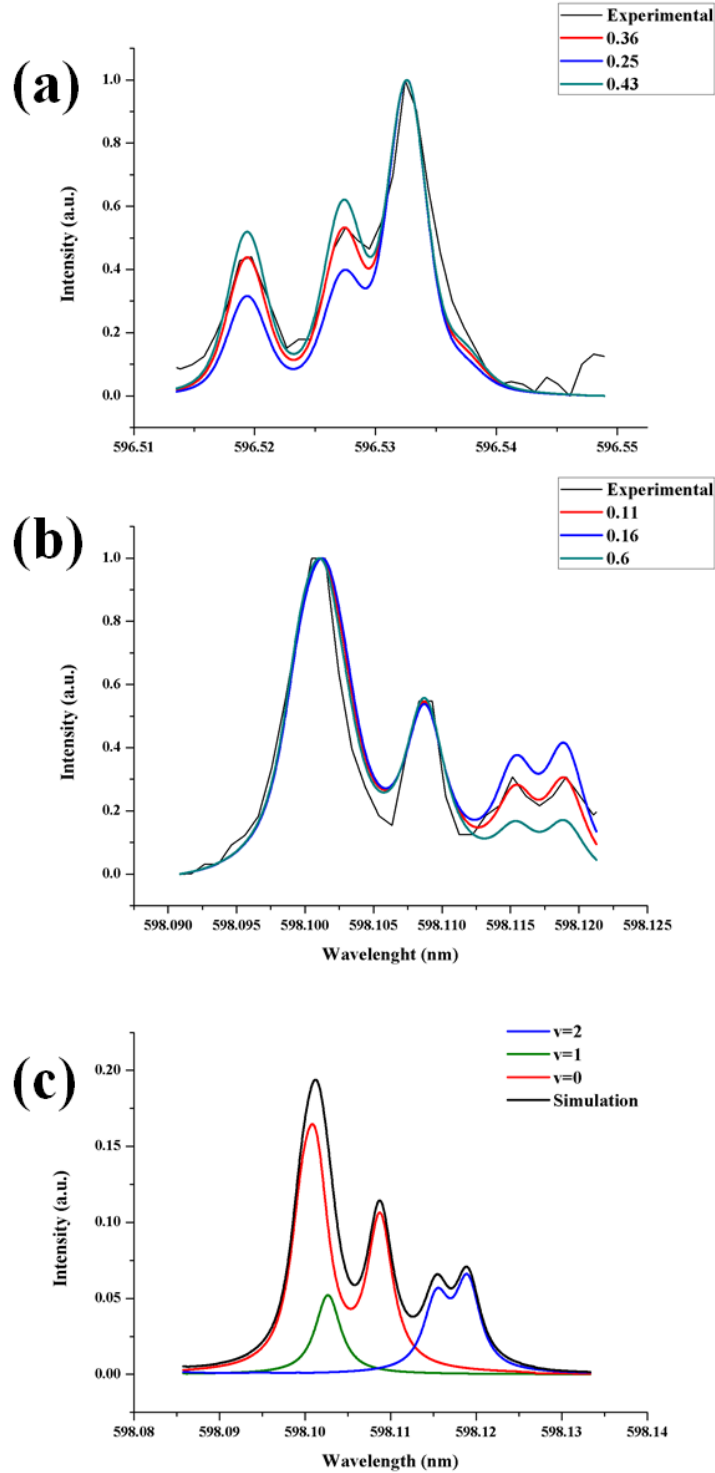


Figure 7: Comparison between experimental and simulated spectra. (a), different population ratio of  $(v'' = 0)/(v'' = 1)$  fixed at 0.36 at 0.044 Torr. (b) different population ratio of  $(v'' = 0)/(v'' = 1)$  fixed at 0.36 and  $(v'' = 0)/(v'' = 2)$  fixed at 0.11 at 0.042 Torr. (c) deconvolution of the simulation with the levels  $(v'' = 0)/(v'' = 1)$  fixed at 0.36 and  $(v'' = 0)/(v'' = 2)$  fixed at 0.11. Both spectra are taking at room temperature.

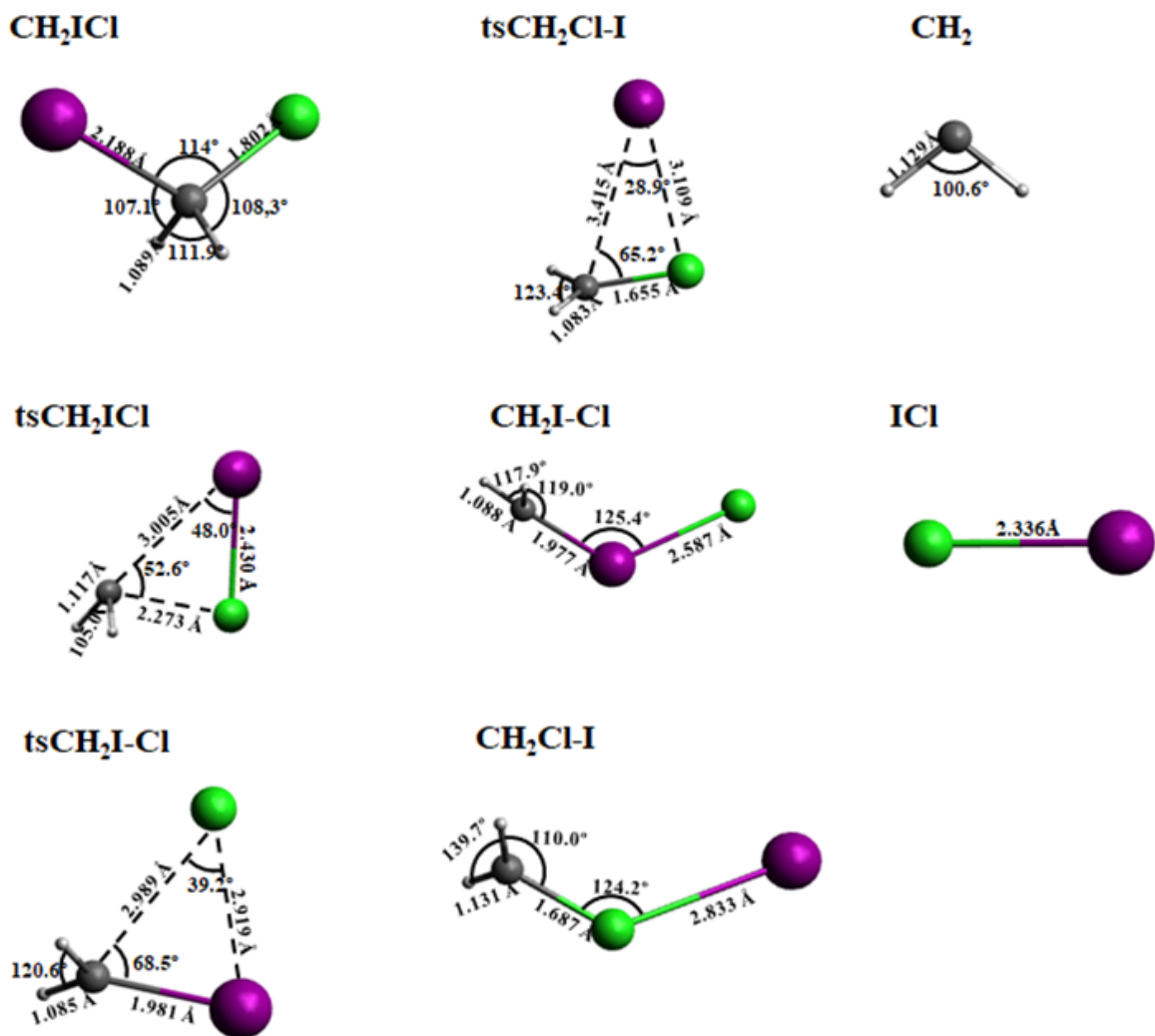


Figure 8: The optimized geometries of reactant, transition state and products for the unimolecular dissociation of  $\text{CH}_2\text{I-Cl}$  using B3LYP/MIDI! level of theory.

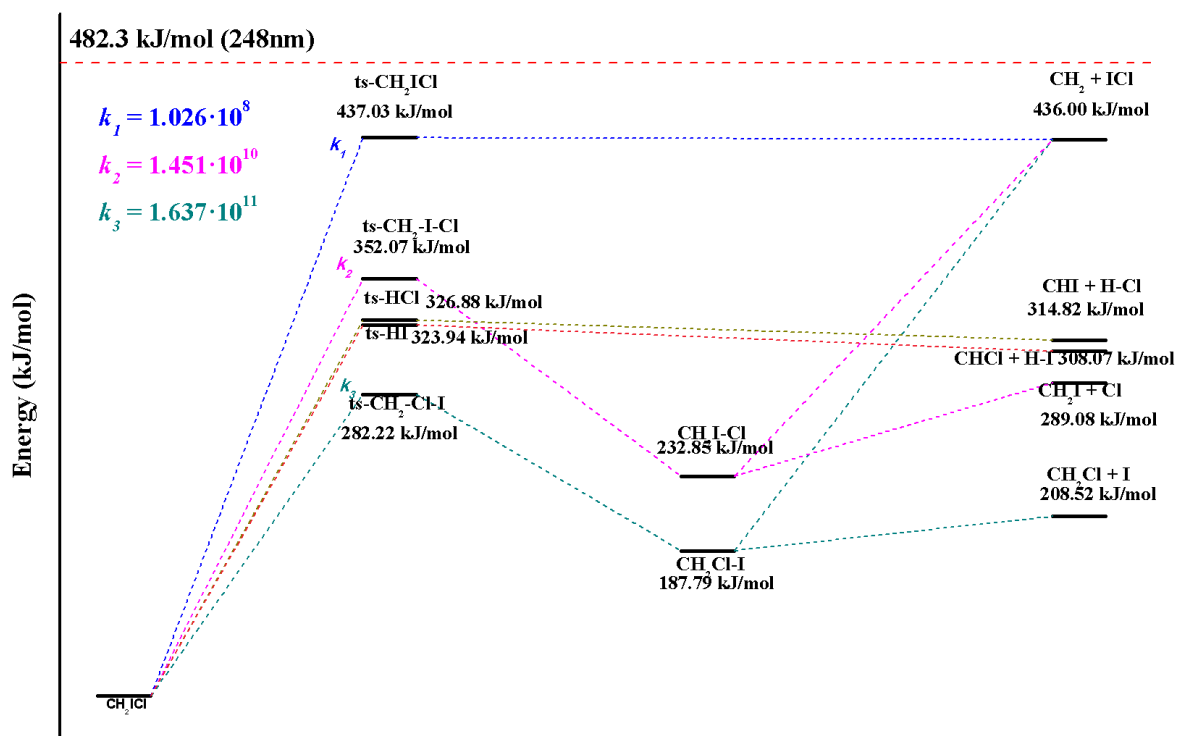


Figure 9: The potential energy surface diagram for the formation of ICl from the CH<sub>2</sub>ICl dissociation channel, as computed using CCSD(T)/MIDI! level of theory with B3LYP/MIDI! zero-point energy corrections. The ICl, I, Cl, HCl and HI dissociation channels of CH<sub>2</sub>ICl, in which the energies are in kJ/mol relative to CH<sub>2</sub>ICl.